

ADVANCED MATERIALS

Supporting Information

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A Ferrotoroidic Candidate with Well-Separated Spin Chains

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Supplementary materials

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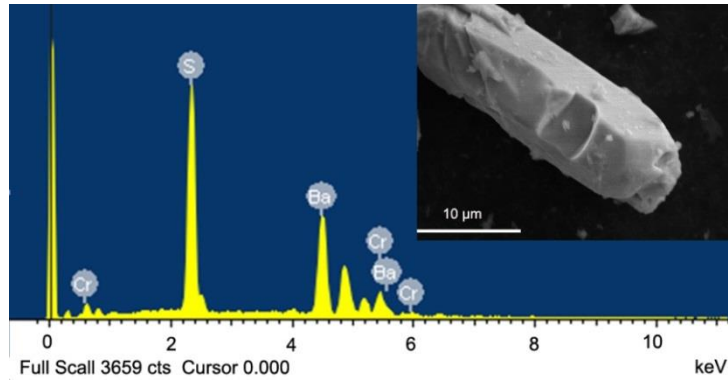


Fig. S1 The energy dispersive X-ray spectrum collected on $\text{Ba}_6\text{Cr}_2\text{S}_{10}$ single crystal. The inset shows the image of the smooth surfaces of the single crystal sample.

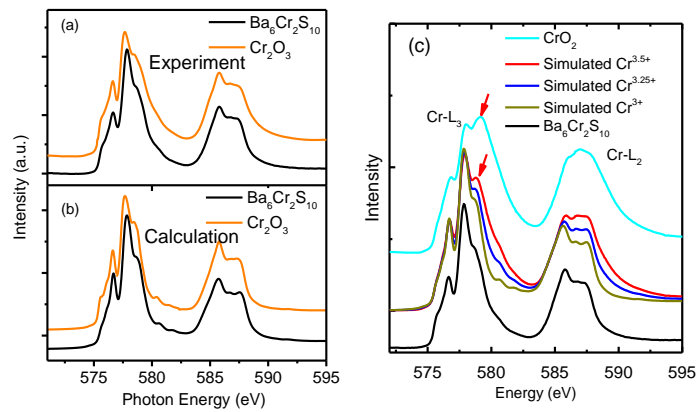


Fig.S2 (a,b) The comparison of experimental and calculated Cr-L_{2,3} XAS spectra of Cr_2O_3 and $\text{Ba}_6\text{Cr}_2\text{S}_{10}$. The experimental spectra for both Cr_2O_3 and $\text{Ba}_6\text{Cr}_2\text{S}_{10}$ can be well reproduced by fully intra-atomic multiplet calculations for a pure Cr^{3+} valence state including ligand field interaction. (c) The comparison of Cr-L_{2,3} XAS spectra of $\text{Ba}_6\text{Cr}_2\text{S}_{10}$, simulated Cr^{3+} , $\text{Cr}^{3.25+}$, $\text{Cr}^{3.5+}$ and CrO_2 with Cr^{4+} . In addition, the comparison of Cr-L_{2,3} XAS of experimental $\text{Ba}_6\text{Cr}_2\text{S}_{10}$ and CrO_2 and the simulated Cr^{3+} , $\text{Cr}^{3.25+}$ and $\text{Cr}^{3.5+}$, firmly excludes the possibility of mixed valence Cr ion in our $\text{Ba}_6\text{Cr}_2\text{S}_{10}$ sample, since a clear peak occurs at ~ 578.9 eV at the Cr-L₃ edge for simulated $\text{Cr}^{3.25+}$ and $\text{Cr}^{3.5+}$ and CrO_2 with Cr^{4+} against a broad weak shoulder in the experimental $\text{Ba}_6\text{Cr}_2\text{S}_{10}$ spectrum (black line).

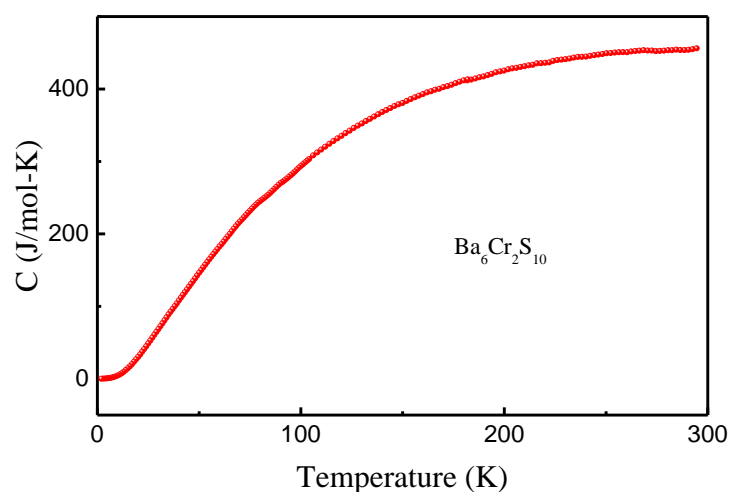


Fig. S3 The specific heat as a function of temperature between 2 K and 300 K measured at zero-field. No abnormality corresponding structural transition is observed, which suggests that the dimerized structure is kept at low temperature. Also, no λ -type abnormality can be found in the specific heat data corresponding to the magnetic transition of ~ 10 K.

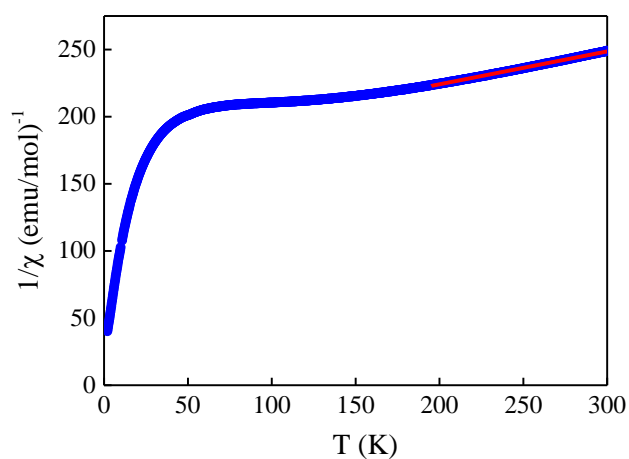


Fig. S4 The inverse of susceptibility as a function of temperature. The red line is the fit by the formula $\chi^{-1} = (T - \theta_p)/C$.

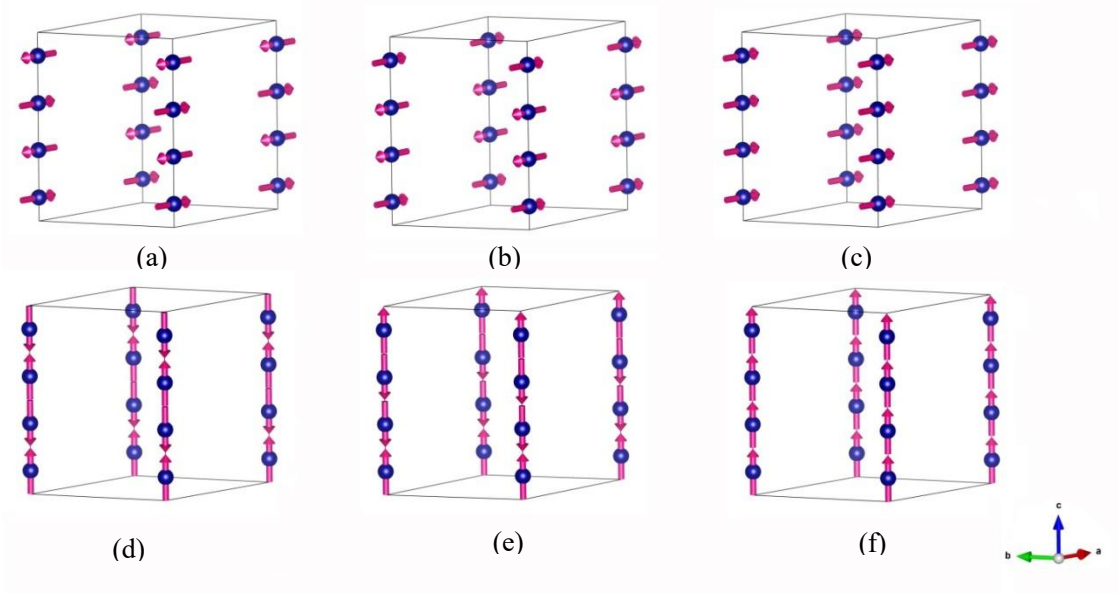


Fig. S5 Different possible magnetic structure modes for $\text{Ba}_6\text{Cr}_2\text{S}_{10}$. (a-c) The spins are oriented along the a axis with the spin configurations of $(\uparrow \downarrow \uparrow \downarrow)$, $(\uparrow \downarrow \downarrow \uparrow)$ and $(\uparrow \uparrow \uparrow \uparrow)$, respectively. (d-e) The same spin configurations with that in (a-c) while the spins are oriented along the c axis. The models of (a-f) are labeled with a -AFM01, a -AFM02, a -FM, c -AFM01, c -AFM02, c -FM, respectively. The three modes of (a), (b) and (f) are allowed by the symmetry of the $Ama'2'$ magnetic space group, but the last two modes are too small to be directly determined from the diffraction experiment.

First-principles calculations were carried out to study the electronic structure and magnetic properties of the system. The Perdew-Burke-Ernzerhof (PBE) exchange-correlation functional has been chosen[1]. The electron correlations associated with $3d$ states of Cr are described by density-functional theory (DFT) + U + SOI (spin-orbit interaction) methods implemented in the Quantum Espresso code[2]. The Vanderbilt ultrasoft pseudopotentials developed for the PBE exchange-correlation functional have been chosen for all the elements throughout all the DFT + U calculations[2]. The values of $U=0$ and 2 eV were used for these calculations. The Monkhorst-Pack sampling of $4 \times 4 \times 3$ has been used for all the calculations. The plane-wave cutoff energies of 544 eV and the threshold of self-consistent-field (SCF) energy convergence of 10^{-5} eV were employed throughout. A linear mix of 10% of the Fock matrix has been used to accelerate the SCF converging process.

All the energies are with reference to the a -AFM01 state with $U=0$ eV. For $U=0$ eV, the

energies of the structure models of (a-f) are 0 eV, 0.254 eV, 0.287 eV, 0.001 eV, 0.230 eV and 0.290 eV, respectively; while for $U=2$ eV, the energies of the models of (a-f) are 5.617 eV, 5.731 eV, 5.698 eV, 5.628 eV, 5.710 eV and 5.700 eV, respectively. From the calculated results, we can see that the *a*-AFM state has the lowest energy along these possible modes, which is consistent with the experimental results.

References:

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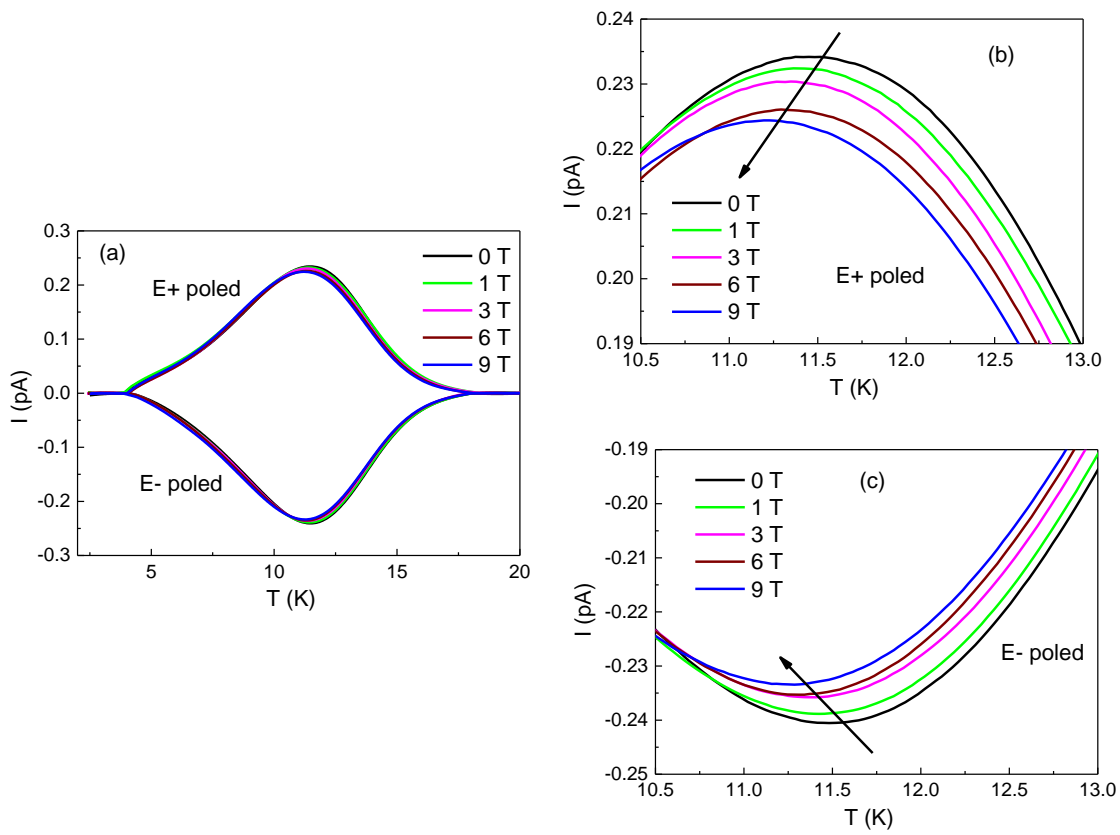


Fig. S6 (a) The pyroelectric current I_p poled from 20 K to 2 K for the polycrystalline sample of $Ba_6Cr_2S_{10}$. The enlarged view for the E+ poled (b) and the E- poled (c), showing the shift of I_p curve toward low temperature under high magnetic field.

Table S1 The positions \vec{r}_i and the magnetic moment \vec{M}_i of Cr cations in Ba₆Cr₂S₁₀. Here, the displacement $\varepsilon=0.00435$ and the magnetic moment $S=1.03(9) \mu_B$, oriented along a axis when ignoring the canted feature. In the main manuscript, the displacement $\delta = \varepsilon \times c = 0.0538(4) \text{ \AA}$. The lattice parameters are $a=b=9.1228(3) \text{ \AA}$ and $c= 12.3643(2) \text{ \AA}$.

Site	r_i^x/a	r_i^y/b	r_i^z/c	m_i^x	m_i^y	m_i^z
Cr	0	0	$7/8+\varepsilon$	S	0	0
Cr	0	0	$5/8-\varepsilon$	S	0	0
Cr	0	0	$3/8+\varepsilon$	S	0	0
Cr	0	0	$1/8-\varepsilon$	S	0	0
Cr	1	0	$7/8+\varepsilon$	S	0	0
Cr	1	0	$5/8-\varepsilon$	S	0	0
Cr	1	0	$3/8+\varepsilon$	S	0	0
Cr	1	0	$1/8-\varepsilon$	S	0	0
Cr	1/2	$\sqrt{3}/2$	$7/8+\varepsilon$	S	0	0
Cr	1/2	$\sqrt{3}/2$	$5/8-\varepsilon$	S	0	0
Cr	1/2	$\sqrt{3}/2$	$3/8+\varepsilon$	S	0	0
Cr	1/2	$\sqrt{3}/2$	$1/8-\varepsilon$	S	0	0
Cr	-1/2	$\sqrt{3}/2$	$7/8+\varepsilon$	S	0	0
Cr	-1/2	$\sqrt{3}/2$	$5/8-\varepsilon$	S	0	0
Cr	-1/2	$\sqrt{3}/2$	$3/8+\varepsilon$	S	0	0
Cr	-1/2	$\sqrt{3}/2$	$1/8-\varepsilon$	S	0	0